

Electrons and holes are modelled using Fermi-Dirac statistics since they are fermions with spin half.

The probability that a particular energy level,  $E$ , is occupied is given by the Fermi-Dirac distribution function

$$F(E) = \frac{1}{1 + e^{\frac{(E-E_F)}{kT}}}$$

where  $E_F$  is the Fermi level energy.

The equation for the current in a diode

$$I = I_0 \left( \exp \left( \frac{eV}{kT} \right) - 1 \right)$$

can be derived by an analysis of the flow of charge carriers across the pn junction based on the distribution of carriers as a function of energy (Fermi-Dirac statistics)

---

$$I = I_0 \left( \exp \left( \frac{eV}{kT} \right) - 1 \right)$$

If the voltage across the diode is zero, then

$$\exp\left(\frac{eV}{kT}\right) = 1$$

and the junction current is zero.

If the voltage across the diode is negative, corresponding to a reverse bias, then the term

$$\exp\left(\frac{eV}{kT}\right) \ll 1$$

and the junction current is  $-I_0$  as expected.

---

When a forward bias is applied, the current increases exponentially.

For  $T = 293\text{K}$

Voltage V	Current A
-10	$-1.0^{-10}$
0	0
0.1	$5.5 \times 10^{-9}$
0.4	$0.9 \times 10^{-3}$
0.5	$4.8 \times 10^{-2}$
0.6	2.65
0.7	144
0.9	$4.3 \times 10^5$

---

A very useful simplification is made if the numerical values

for Boltzman's constant,  $1.38 \times 10^{-23} JK^{-1}$ ,

the electronic charge,  $1.6 \times 10^{-19} C$  and

the assumed room temperature of  $t = 20^{\circ}C = 293K$

are substituted to give:—

$$\frac{kT}{e} = \frac{1.38 \times 10^{-23} \times 293}{1.6 \times 10^{-19}} = 0.025V = 25mV$$

---

The junction diode forward current then becomes:—

$$I = I_0 \exp \left( \frac{V}{25mV} \right)$$

This is a result which we will use very frequently in our analysis of circuits.

---